



Theory of Solar Radar Experiments: Combination Scattering by Anisotropic Langmuir Turbulence

Licentiate seminar
by

Mykola Khotyaintsev

Dept. of Astronomy and Space Physics,
Uppsala University, Sweden

November 8, 2005. Uppsala, Sweden



Presentation outline

1. The Sun

- Overview
- Prominence, flares, and coronal mass ejections (CME's)
- Type III solar radio bursts

2. Solar radar experiments

- Overview
- Main experimental results

3. Theory of radar reflections from the Sun

- Specular reflection
- Volume scattering by density fluctuations
- Induced (combination) scattering by wave turbulence

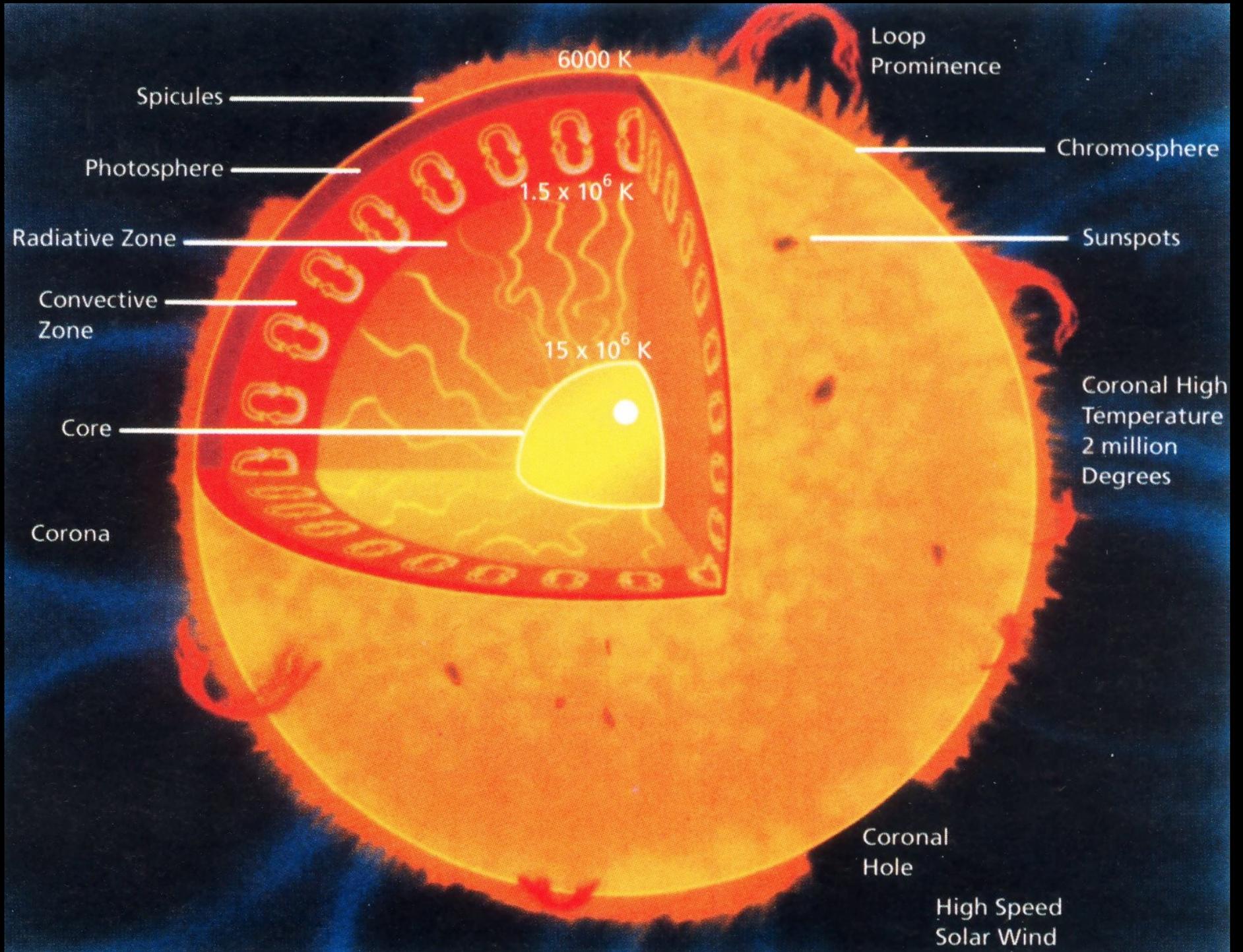
4. Paper

M.V. Khotyaintsev, V.N. Melnyk, Bo Thide' and O.O. Konovalenko
“Combination scattering by anisotropic Langmuir turbulence with
application to solar radar experiments”, *Solar Physics*, in press, 2005

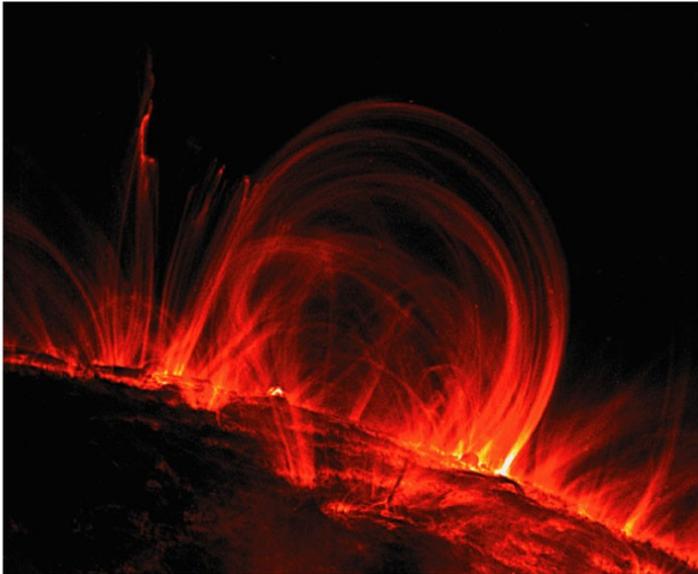
Motivation of solar radar experiments

- radar is an additional tool for solar study
- to investigate dynamics of the solar corona
- solar wind acceleration
- coronal electron density profile
- magnetic field probing
- remote sensing of plasma wave turbulence
- detection of coronal mass ejections (CMEs)
- recent experiment proposals by Thide' (2002), Coles (2004), Rodriguez (2004)

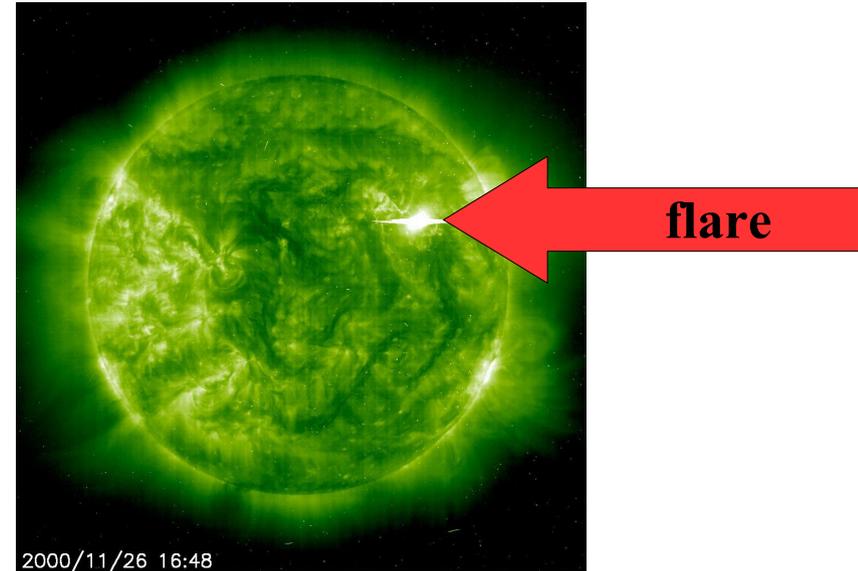
The structure of the Sun



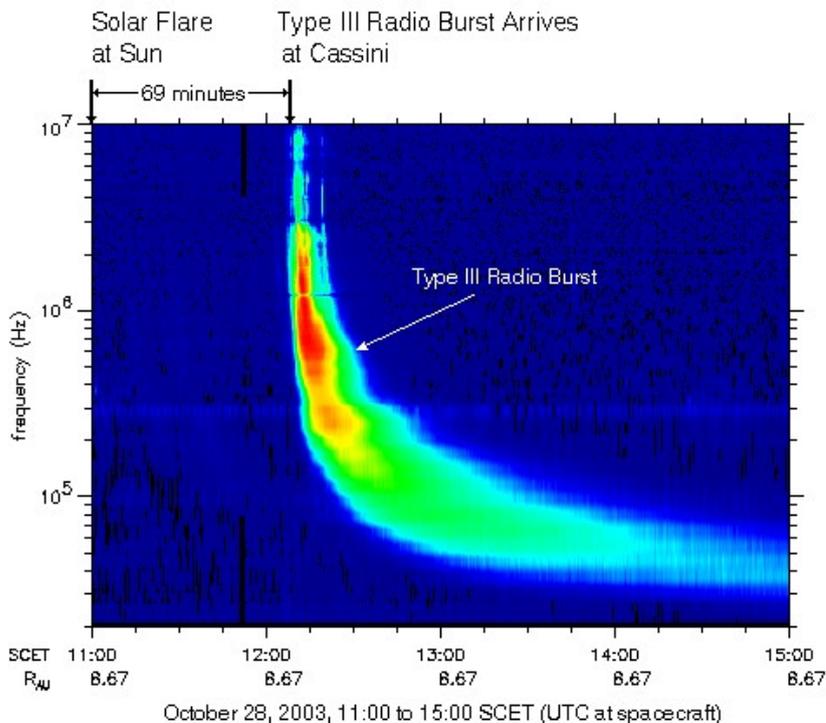
Prominences, flares and type III radio bursts



Copyright © 2004 Pearson Education, publishing as Addison Wesley.



X-ray observation with the EIT on SOHO



Type III solar burst

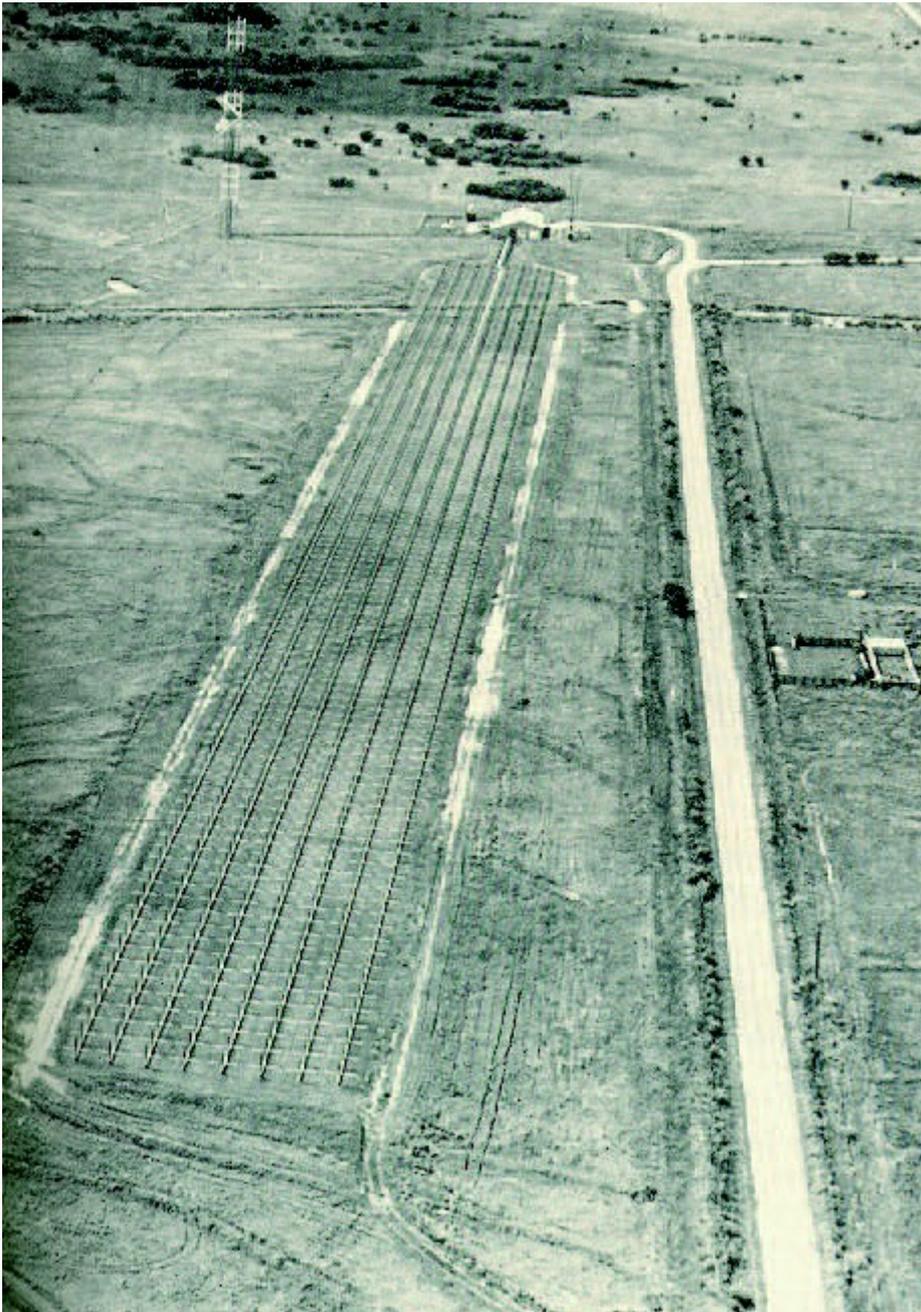
- are associated with solar flares
- attributed to beams of electrons with velocities $v = 0.2 - 0.6 c$
- generate Langmuir turbulence
- radiation is emitted at the local plasma frequency \Rightarrow density probe

Solar radar experiments

RADAR = RAdio Detection And Ranging

- 1940s – first radar studies of space objects (the moon)
- 1959 – first solar radar experiment at 25 MHz by a Stanford group
- 1961-1969 – daily experiments at 38 MHz by an MIT group at El Campo, Texas
- 1977 and 1978 an attempt to observe scatter from Langmuir waves in the corona using a **2380 MHz**, transmitter at Arecibo
no echo observed most likely because of the **too high frequency** used
- 1996-1998 – experiment at **9 MHz** at the Russian Sura transmitter and the Ukrainian UTR-2 radio telescope
no echo observed because of the **too low frequency** used

The El Campo solar radar experiment



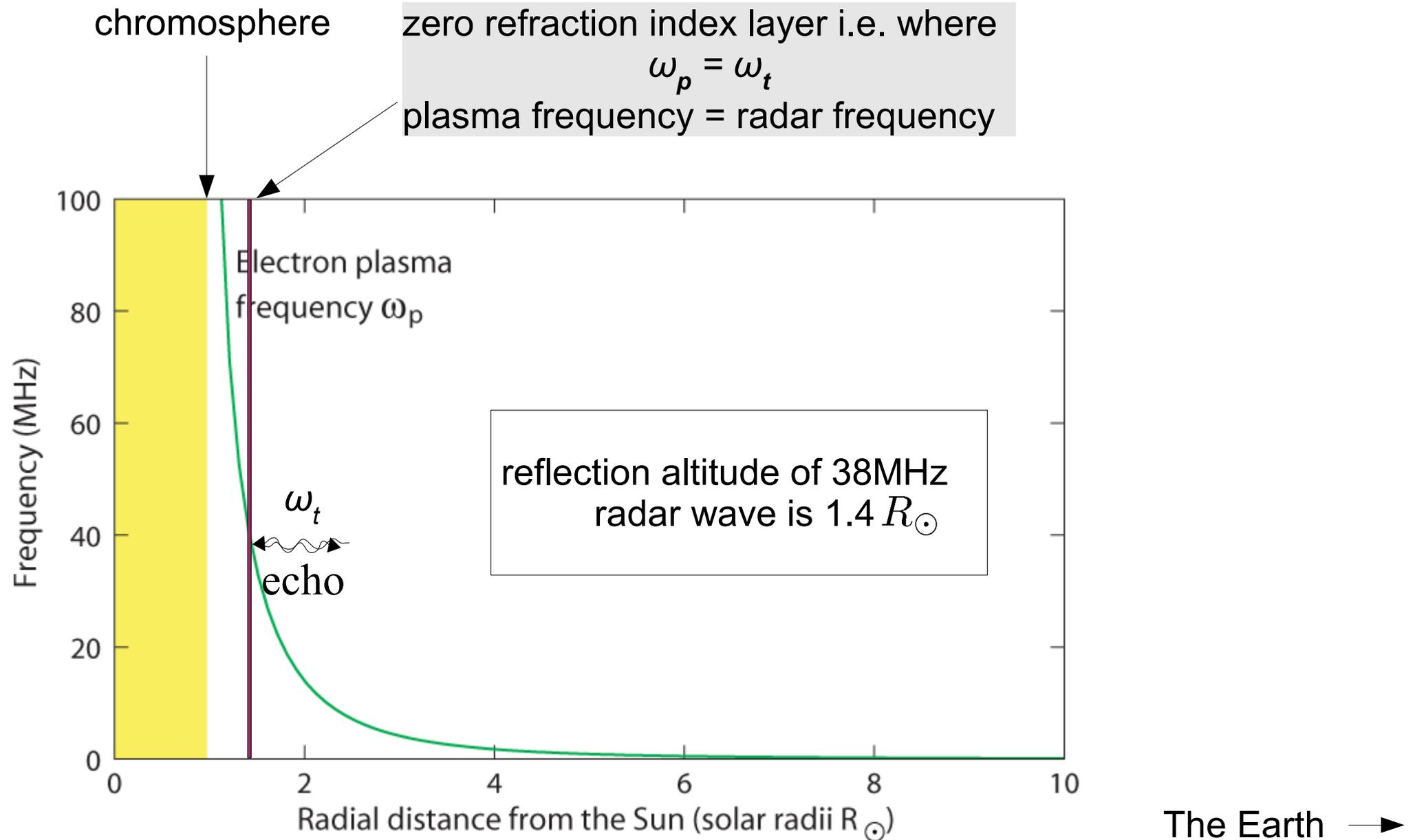
Operating frequency:	38.25 MHz
Total power:	500 kW
Beam size:	$1^\circ \times 6^\circ$
Size of the Sun:	0.5°
Gain:	32-36 dB
Eff. radiated power:	1300 MW

Operation mode:

16 min. of transmission followed by
16 min. of reception

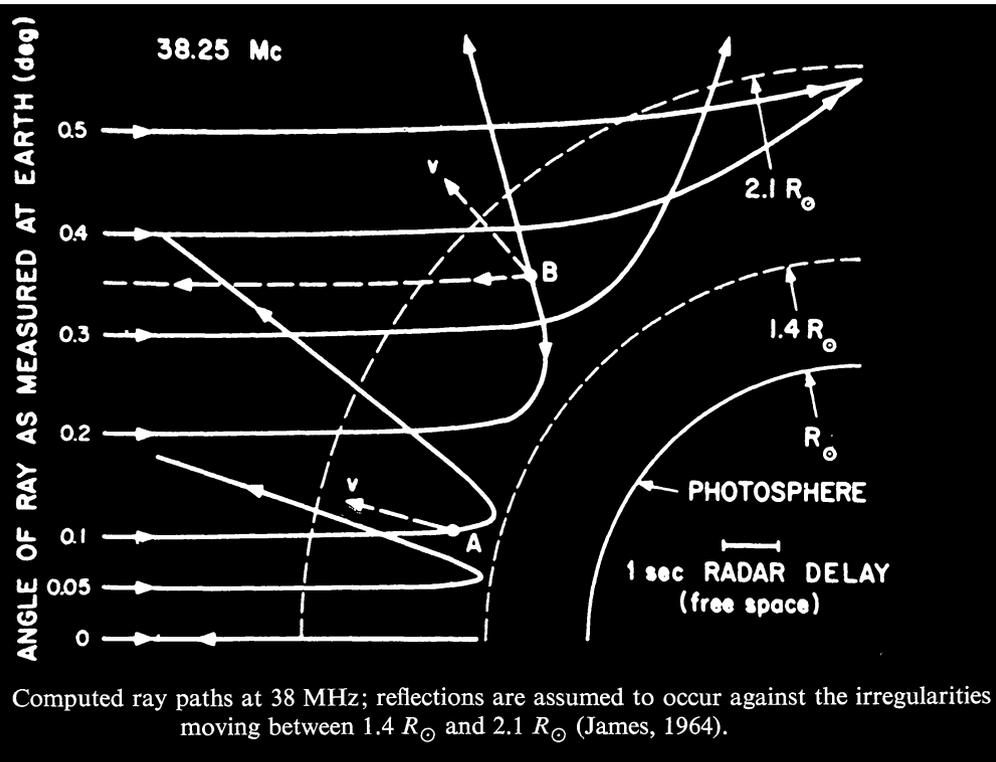
light travels from the Sun
to the Earth in 8 min

Theory: specular reflection of a radio wave the corona (1D)

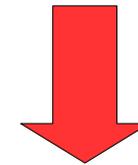


Specular reflection from the spherically-symmetric corona

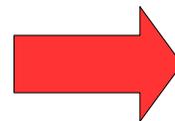
Ray paths at 38 MHz in the corona



- reflection occurs from the layer with the zero refracting index, i.e. where
$$\omega_p = \omega_t$$
plasma frequency = radar frequency
- reflection from a **rough** sphere
- collisional absorption



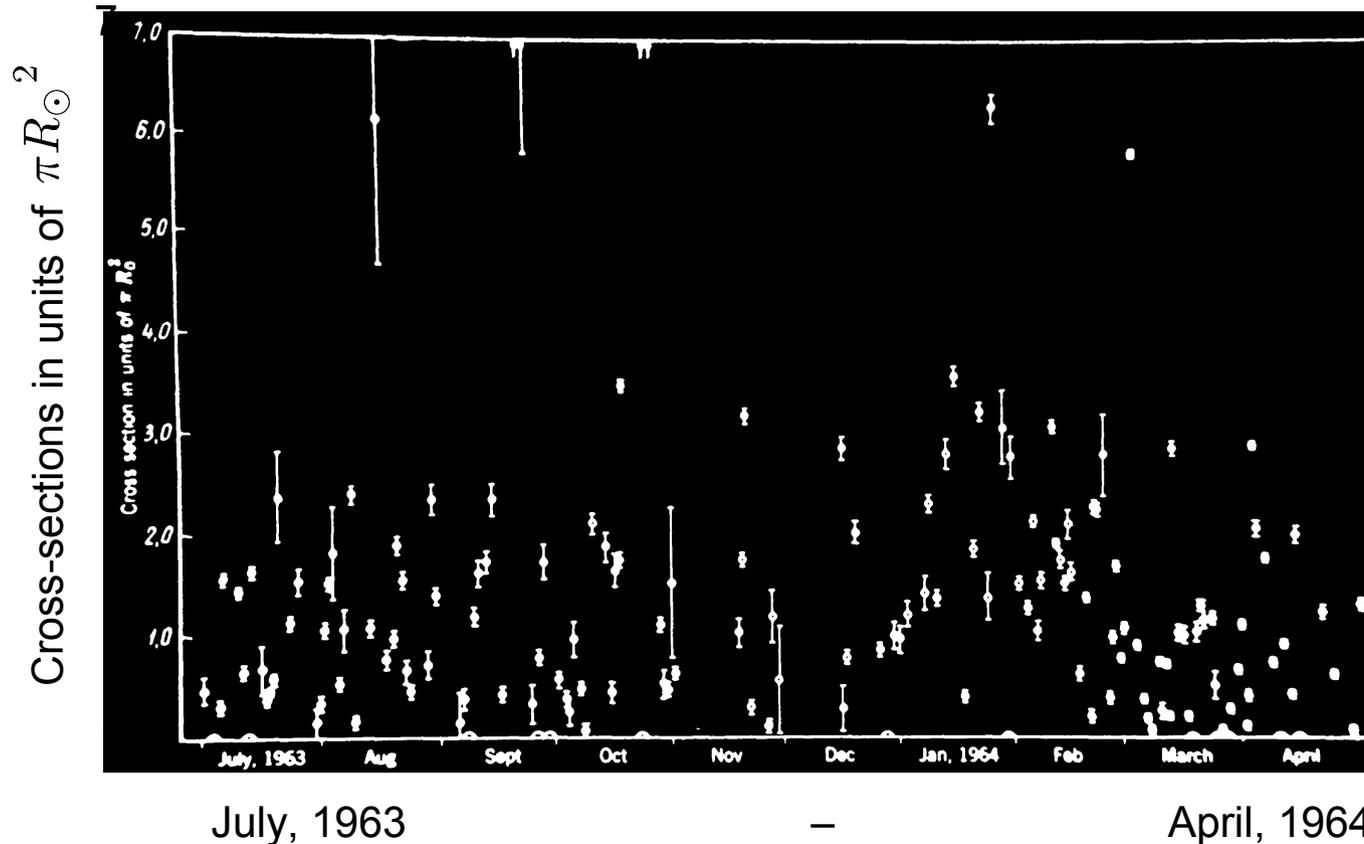
Volume scattering by large-scale ($L \gg \lambda$) density irregularities



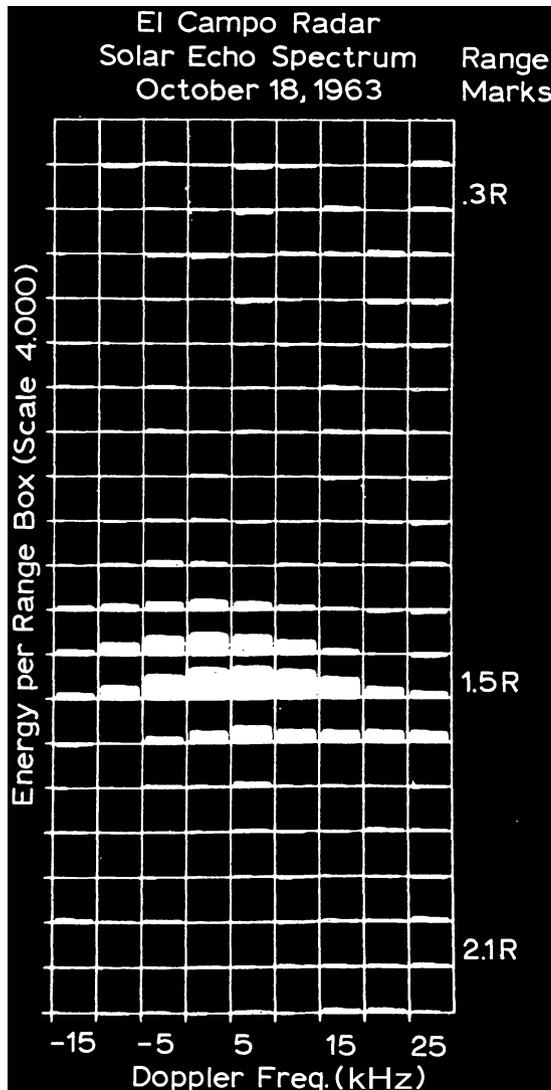
Cross-section $S \approx 1.5 \pi R_{\odot}^2$

Solar cross-sections detected with the El Campo radar

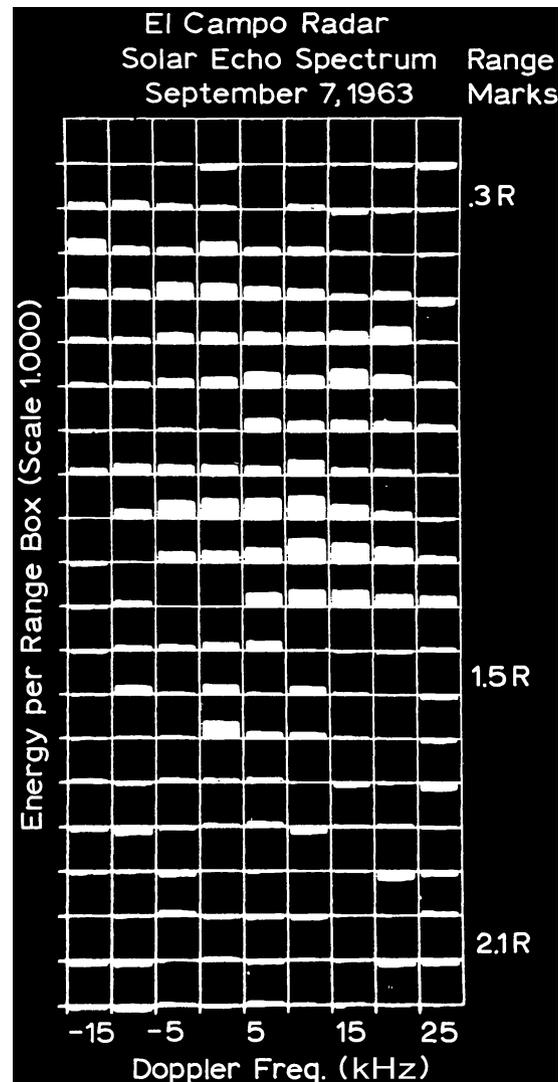
- Cross-sections observed are in the range $0 < S < 800 \pi R_{\odot}^2$
- The majority of cross-section $S = 0 - 4 \pi R_{\odot}^2$



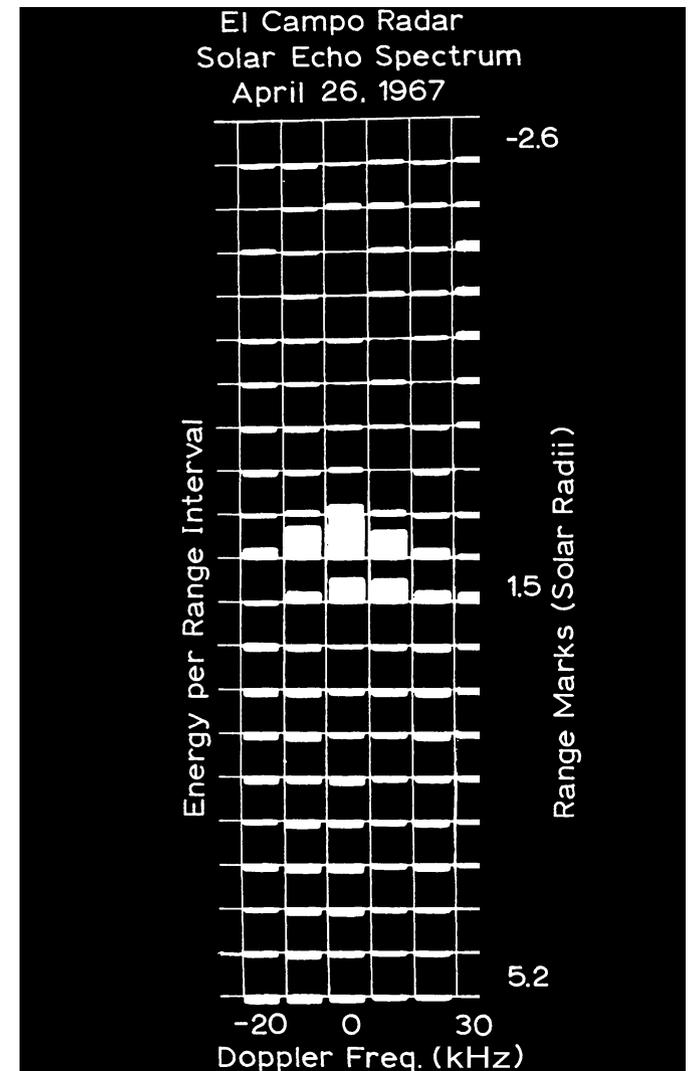
El Campo radar echo spectrum



Type A (70%)



Type B: echo is coming from different altitudes

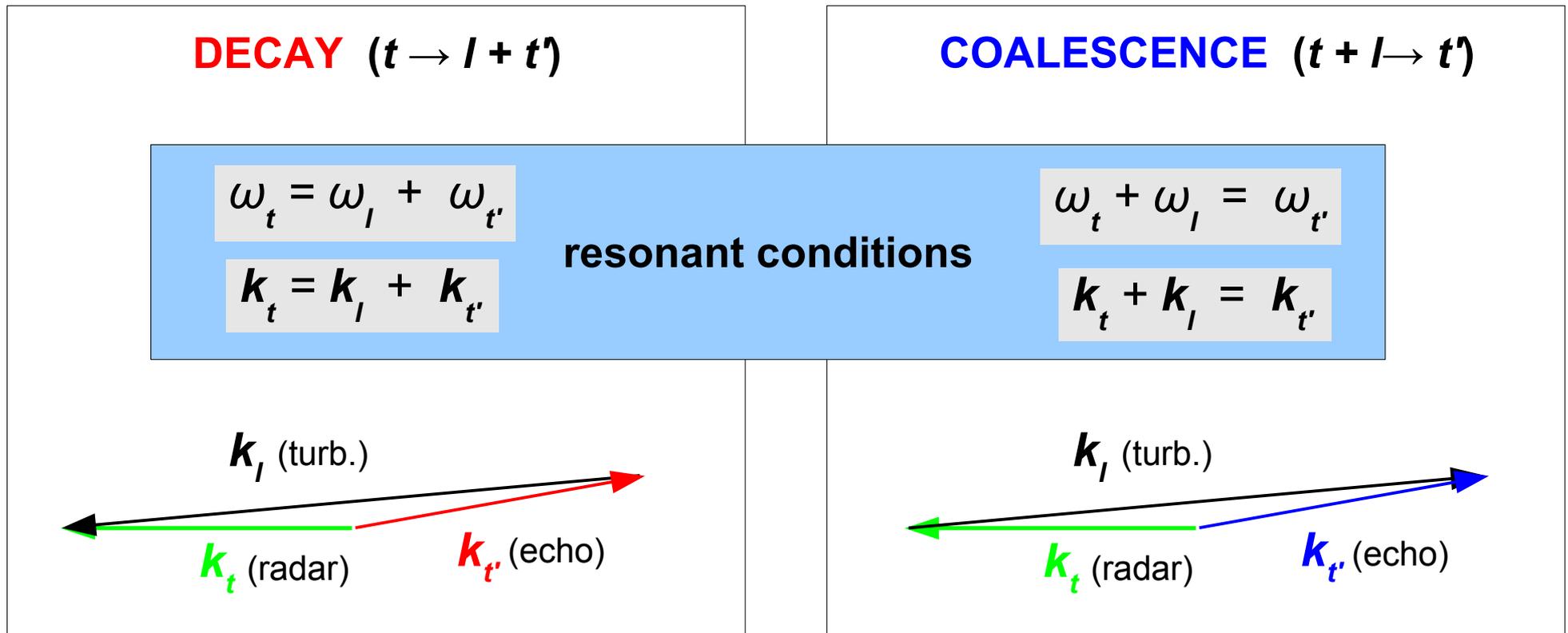


Extreme case of type B spectrum

- experimental results could not be explained by the specular reflection theory

Induced (combination) scattering by a wave turbulence

- corona contains areas of localized plasma wave turbulence:
Langmuir (l) and ion-sound (s)
- echo signal may be formed due to resonant interaction of the radar wave with the waves of the turbulence (induced scattering)



- radar wave ω_t is scattered into two satellites:

$$\omega_t - \omega_l \text{ (red-shifted) and } \omega_t + \omega_l \text{ (blue-shifted)}$$

- qualitative description is given by **the kinetic wave equation** (Tsytovich, 1970)

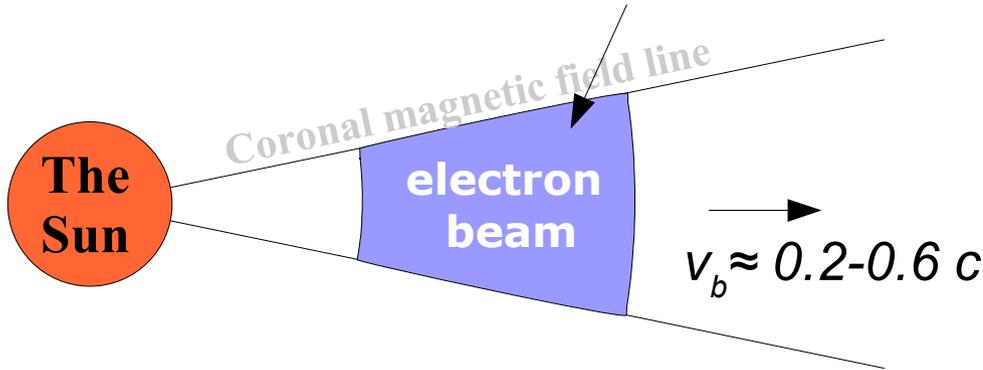
Paper : Combination scattering by anisotropic Langmuir turbulence with application to solar radar experiments

Motivation

- can radar echoes come from type III burst turbulence?
 - properties of the echo and what can we derive from them?
 - to give hints for future radar experiments: transmitting freq., receiving bandwidth, etc.
-
- We assume existence of a Langmuir turbulence generated by type III electrons (I)
 - Existence of the radar wave (t)
 - Induced scattering process in a weak turbulence limit
 - Focus on backscattering

What is a type III solar burst?

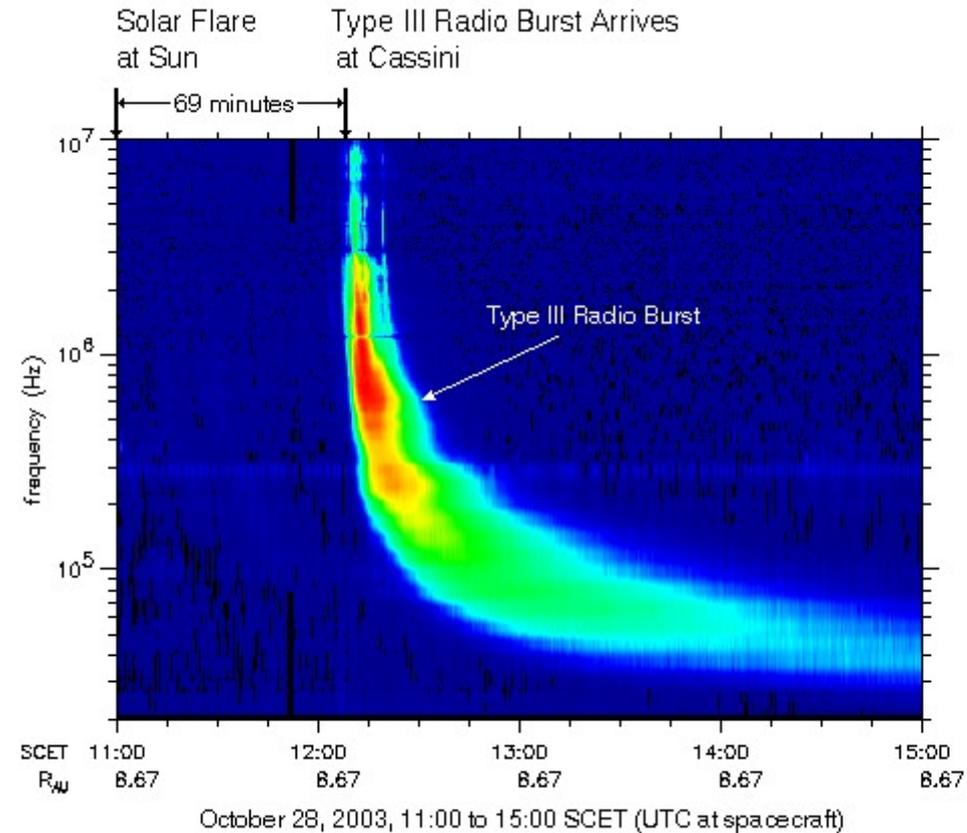
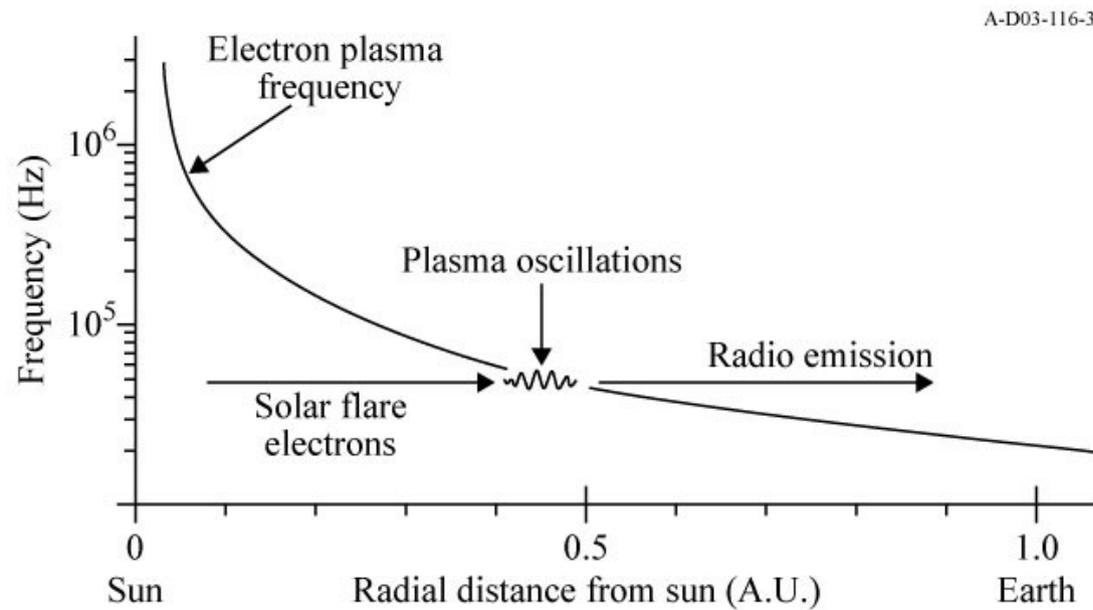
Beam instability => Langmuir turbulence



Radiation generation:

- 1) $l \rightarrow l' + s$ (s – ion-sound)
- 2) $l + s \rightarrow t$

t' (echo) \rightarrow $\omega_l \gg \omega_s \Rightarrow \omega_t \approx \omega_p$ (radiation at plasma frequency) \leftarrow t (radar)



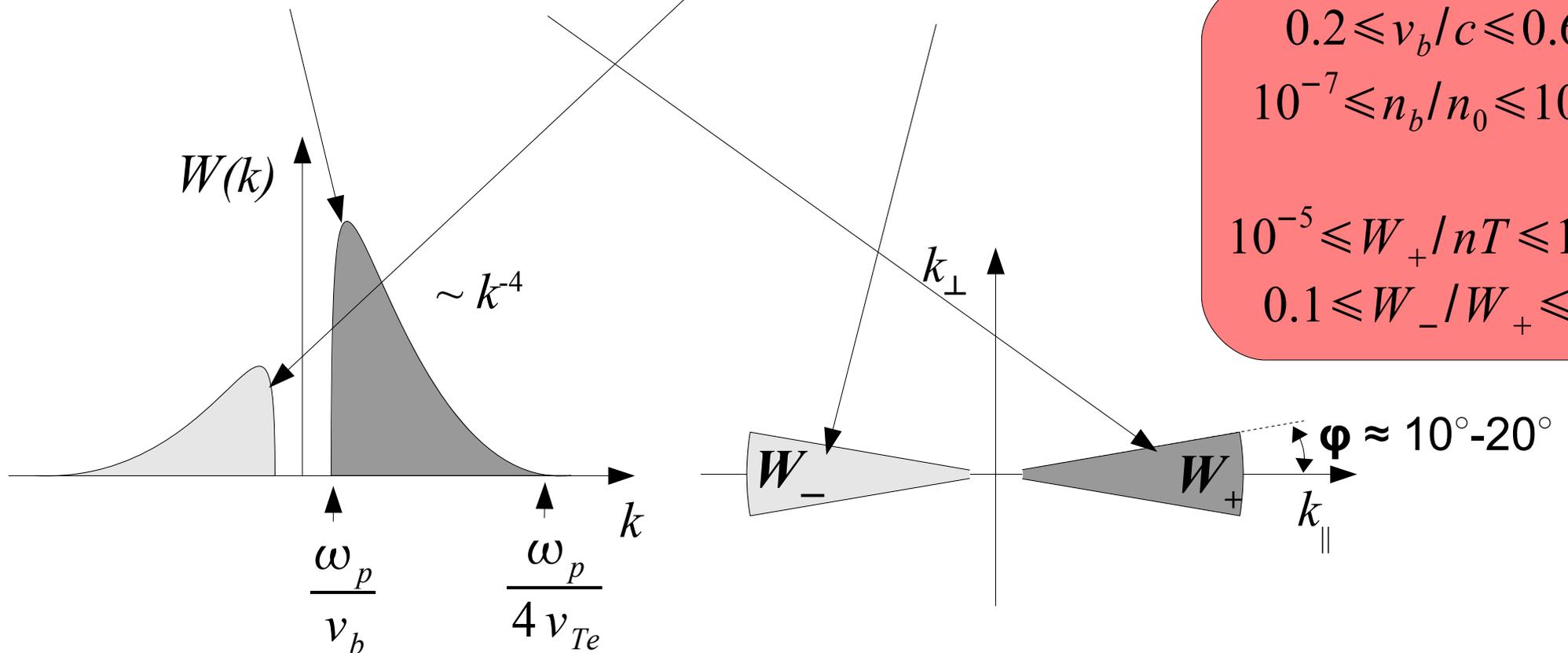
Langmuir turbulence spectral energy density $W(k)$

Mel'nik et. al. (1999)

Primary turbulence W_+
(due to beam instability)

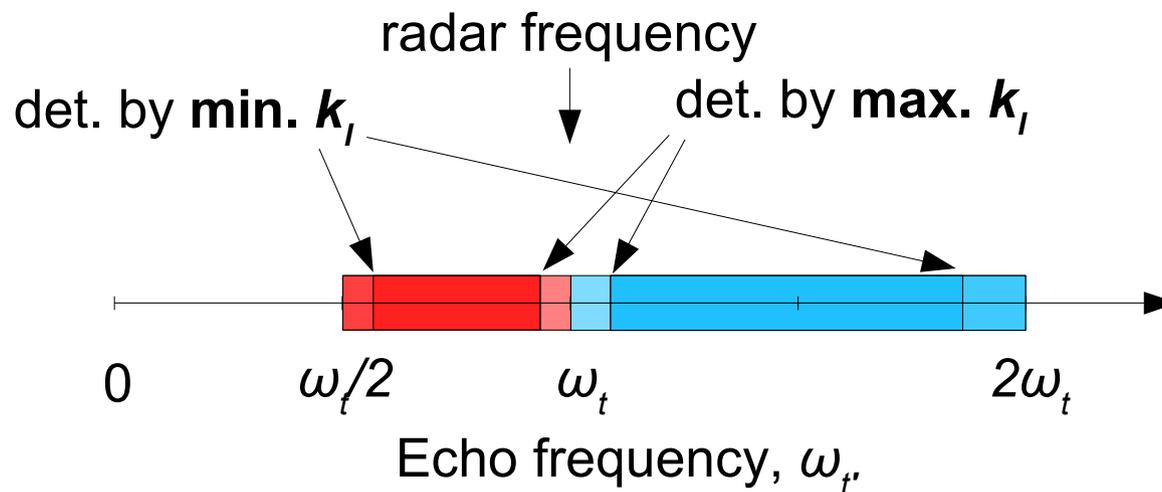
Secondary turbulence W_-
(due to $l \rightarrow l' + s$)

$0.2 \leq v_b/c \leq 0.6$
 $10^{-7} \leq n_b/n_0 \leq 10^{-5}$
 $10^{-5} \leq W_+/nT \leq 10^{-3}$
 $0.1 \leq W_-/W_+ \leq 1$



Echo frequency

- echo frequency shift $|\Delta\omega| = |\omega_{t'} - \omega_t| \approx \omega_p$ (plasma frequency)
- $\Delta\omega < 0$ (decay) $\Delta\omega > 0$ (coalescence)
- echo frequencies lie in a limited range



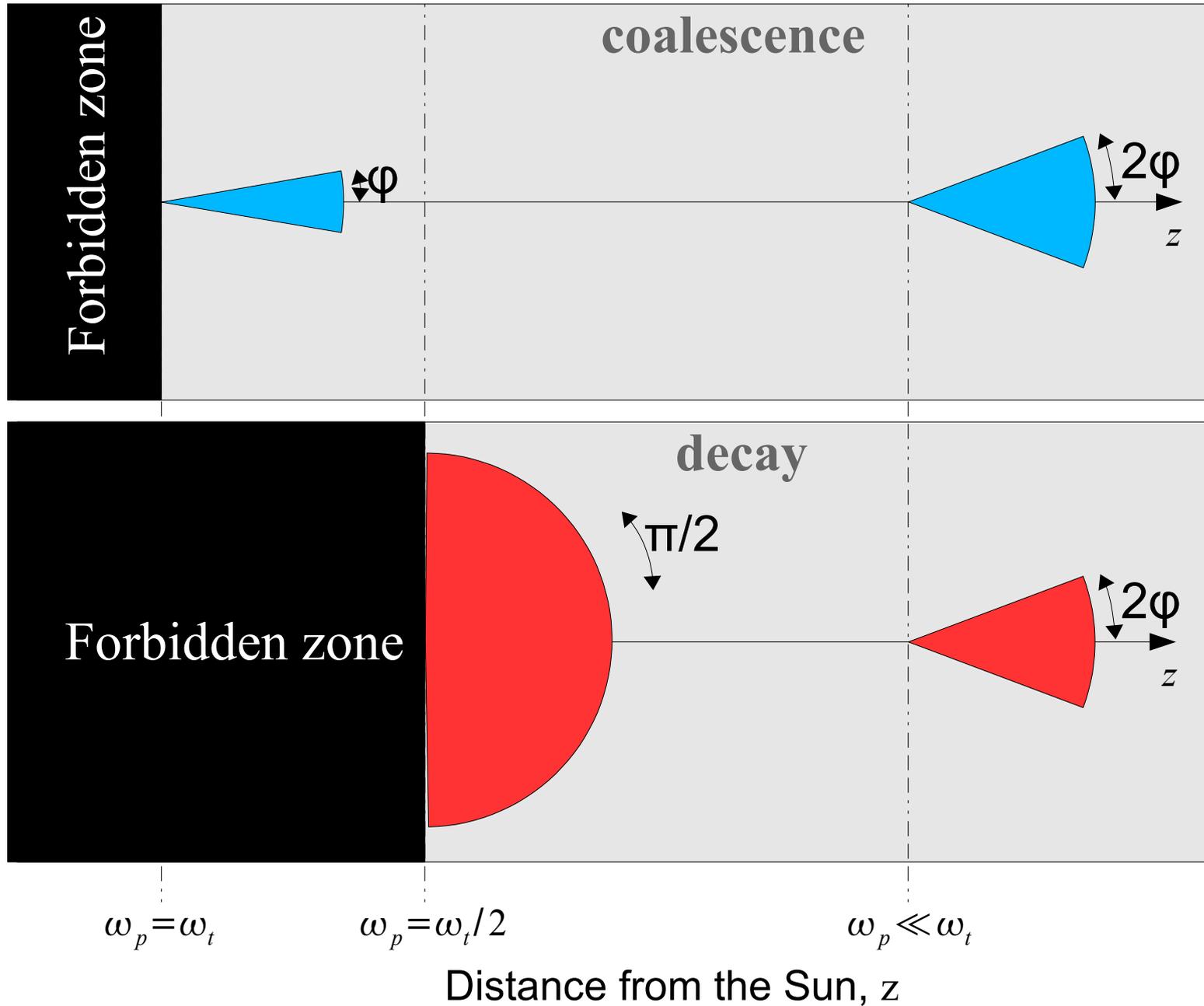
minimal $|\Delta\omega|$ is determined by **max. k_i** of turbulence spectrum

maximal $|\Delta\omega|$ is determined by **min. k_i** (beam velocity)

- echoes with $\Delta\omega > 0$ are due to scattering by the **primary** turbulence \mathbf{W}_+
- echoes with $\Delta\omega < 0$ are due to scattering by the **secondary** turbulence \mathbf{W}_-
- turbulence **anisotropy** \Rightarrow scattering by $\Delta k \ll k$

Echo angular spread

φ - angular spread of Langmuir waves



Efficiency of the scattering process

- The kinetic wave equation (*Tsytovich, 1970*) for the intensity of the incident radar wave \mathbf{W}_{kt}^t may be simplified to the radiation transfer equation

$$\frac{dW_{k_t}^t}{dz} = -[\mu_+(z) + \mu_-(z)]W_{k_t}^t$$

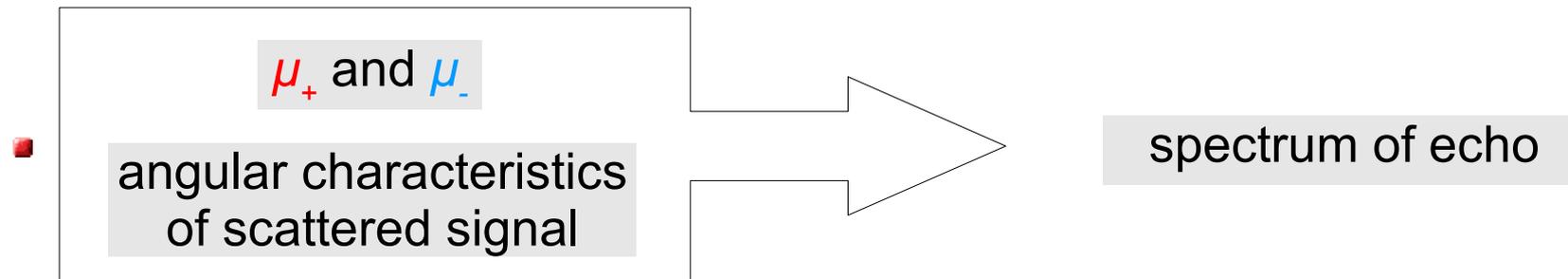
μ_+ and μ_- are coeff. of attenuation of the radar wave due to scattering

- obtained analytical expressions for μ_+ and μ_- show:

decay is most efficient at $\omega_p \approx \omega_t/2$

coal. is most efficient at $\omega_p \approx \omega_t$

- the scattering process is “optically thick” ($\tau = \int_{z_1}^{z_2} \mu_+ + \mu_- dz \geq 1$) for $W \geq 10^{-5} n\kappa T$



Summary of the paper

- Studied the process of scattering of a radar beam by an anisotropic Langmuir turbulence
- Showed, that the frequency of radar echo is within a limited frequency range
- Angular spread of blue-shifted and red-shifted echo differ dramatically
- Decay and coalescence are most efficient at alt. of $\omega_p \approx \omega_t/2$ and ω_t
- Obtained estimates of an echo spectrum
- Minimum turbulence level needed for a reflection is $W \approx 10^{-5} nT$
- Radar experiments may be used to study the spectrum of the beam-generated Langmuir turbulence

Summary and outlook

- Solar radar can be a useful tool for solar studies

- There is a need in new radar experiments

- There is a need in further development of theory

- Future work: theory of radar reflections from CMEs